

AD-A167 600

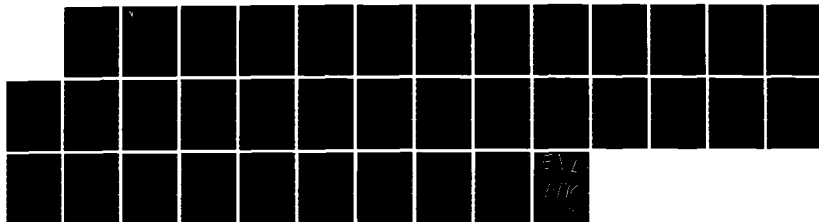
ANNIHILATION OF ANTIPROTONS IN HEAVY NUCLEI(U) LAWRENCE  
LIVERMORE NATIONAL LAB CA D L MORGAN APR 86 UCID-20724  
AFRPL-TR-86-011 MIPR-RL-59004

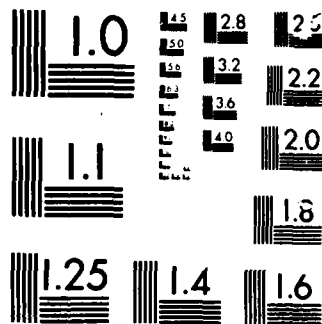
1/1

UNCLASSIFIED

F/G 20/8

NL





MICROCOPY

CHART

AD-A167 608



AFRPL TR-86-011

AD:

(12)

al Report  
the period  
April 1985 to  
December 1985

## Annihilation of Antiprotons in Heavy Nuclei

April 1986

Author:  
D. L. Morgan, Jr.

Lawrence Livermore National Laboratory  
7000 East Avenue  
P.O. Box 808  
Livermore, CA 94550

UCID-20724  
MIPR: RPL 59004

### Approved for Public Release

Distribution is unlimited. The AFRPL Technical Services Office has reviewed this report, and it is releasable to the National Technical Information Service, where it will be available to the general public, including foreign nationals.

RECEIVED  
MAY 22 1986

*prepared for the:*

**Air Force  
Rocket Propulsion  
Laboratory**

Air Force Space Technology Center  
Space Division, Air Force Systems Command  
Edwards Air Force Base,  
California 93523-5000

OTIC FILE COPY

86 5 21 031

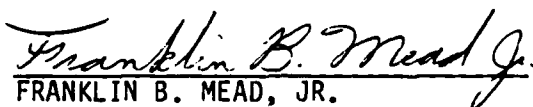
## NOTICE

When U.S. Government drawings, specifications, or other data are used for any purpose other than a definitely related government procurement operation, the government thereby incurs no responsibility nor any obligation whatsoever, and the fact that the government may have formulated, furnished, or in any way supplied the said drawings, specifications, or other data, is not to be regarded by implication or otherwise, or conveying any rights or permission to manufacture, use, or sell any patented invention that may in any way be related thereto.


## FOREWORD


This report was prepared for the Air Force Rocket Propulsion Laboratory (AFRPL) under MIPR RPL-59004 by the Lawrence Livermore National Laboratory (LLNL) under the auspices of the U. S. Department of Energy. The work was performed at the LLNL during the period 1 April 1985 to 31 December 1985. Principal investigator for the LLNL was Mr David L. Morgan, Jr. Project Manager for the AFRPL was Dr Franklin B. Mead, Jr.

This technical report has been reviewed and is approved for publication and distribution in accordance with the distribution statement on the cover and on the DD Form 1473.

  
FRANKLIN B. MEAD, JR.  
Project Manager

FOR THE DIRECTOR

  
EDWARD S. HUSTON, Maj, USAF  
Chief, Space Propulsion Branch

  
EDWARD L. IBBOTSON, Lt Col, USAF  
Deputy Chief, Liquid Rocket Division

• UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE

| REPORT DOCUMENTATION PAGE  |   |   |                                 |          |               |
|--|---|---|---------------------------------|----------|---------------|
| 1a. REPORT SECURITY CLASSIFICATION<br>UNCLASSIFIED   |   | 1b. RESTRICTIVE MARKINGS  |                                 |          |               |
| 2a. SECURITY CLASSIFICATION AUTHORITY  |   | 3. DISTRIBUTION/AVAILABILITY OF REPORT<br>Approved for Public Release: Distribution is Unlimited. |                                 |          |               |
| 2b. DECLASSIFICATION/DOWNGRADING SCHEDULE  |   |   |                                 |          |               |
| 4. PERFORMING ORGANIZATION REPORT NUMBER(S)<br>UCID-20724  |   | 5. MONITORING ORGANIZATION REPORT NUMBER(S)<br>AFRPL-TR-86-011                                    |                                 |          |               |
| 6a. NAME OF PERFORMING ORGANIZATION<br>Lawrence Livermore National Laboratory  | 6b. OFFICE SYMBOL<br>(If applicable)          | 7a. NAME OF MONITORING ORGANIZATION<br>Air Force Rocket Propulsion Laboratory                     |                                 |          |               |
| 6c. ADDRESS (City, State and ZIP Code)<br>7000 East Avenue<br>P.O. Box 808<br>Livermore, CA 94550  |   | 7b. ADDRESS (City, State and ZIP Code)<br>AFRPL/LKC<br>Edwards AFB, CA 93523-5000                 |                                 |          |               |
| 8a. NAME OF FUNDING/SPONSORING ORGANIZATION  | 8b. OFFICE SYMBOL<br>(If applicable)          | 9. PROCUREMENT INSTRUMENT IDENTIFICATION NUMBER<br>MIPR- RPL-59004                                |                                 |          |               |
| 8c. ADDRESS (City, State and ZIP Code)   |   | 10. SOURCE OF FUNDING NOS.  |                                 |          |               |
|  |   | PROGRAM ELEMENT NO.   | PROJECT NO.                     | TASK NO. | WORK UNIT NO. |
| 11. TITLE (Include Security Classification)<br>Annihilation of Antiprotons in Heavy Nuclei (U)   |   | 62302F  | 5730                            | 00       | IV            |
| 12. PERSONAL AUTHOR(S)<br>Morgan, David L., Jr.  |   |   |                                 |          |               |
| 13a. TYPE OF REPORT<br>Final   | 13b. TIME COVERED<br>FROM 85/04/1 TO 85/12/31 | 14. DATE OF REPORT (Yr., Mo., Day)<br>86/4  | 15. PAGE COUNT<br>35            |          |               |
| 16. SUPPLEMENTARY NOTATION   |   |   |                                 |          |               |
| 17. COSATI CODES   |   | 18. SUBJECT TERMS (Continue on reverse if necessary and identify by block number)                 |                                 |          |               |
| FIELD  | GROUP   | SUB. GR.  |                                 |          |               |
| 20   | 08  |   |                                 |          |               |
|  |   | Antiprotons, Heavy nuclei, Annihilation, Space propulsion, Annihilation medium, Working fluid     |                                 |          |               |
| 19. ABSTRACT (Continue on reverse if necessary and identify by block number)   |   |   |                                 |          |               |
| <p>A literature survey was conducted to find information relevant to the annihilation of a low-energy (sub-MeV) antiproton in a heavy nucleus. Such information is important to the use of nuclear fragments from the annihilation to heat a working fluid for space propulsion. The particular piece of information desired was the fraction of annihilation energy that becomes the kinetic energy of charged nuclear fragments emitted after the annihilation. The experimental and theoretical information located was sufficient to allow calculation of that energy fraction. Its value is about 10% for nuclei as heavy as silicon or greater and 20% for very heavy nuclei when the energy of fission fragments is included. Both values are less than the fraction of annihilation energy (38%) that becomes the kinetic energy of charged pions from the annihilation of an antiproton with a proton (hydrogen nucleus). These values are relevant to the choice of a working fluid that absorbs a portion of the annihilation energy and forms the exhaust for a rocket powered by antiproton annihilation. Although it is easier to couple the energy of charged nuclear fragments to the working fluid,</p> |   |   |                                 |          |               |
| 20. DISTRIBUTION/AVAILABILITY OF ABSTRACT<br>UNCLASSIFIED/UNLIMITED <input checked="" type="checkbox"/> SAME AS RPT <input type="checkbox"/> DTIC USERS <input type="checkbox"/>   |   | 21. ABSTRACT SECURITY CLASSIFICATION (over)<br>UNCLASSIFIED                                       |                                 |          |               |
| 22a. NAME OF RESPONSIBLE INDIVIDUAL<br>FRANK B. MEAD, JR.  |   | 22b. TELEPHONE NUMBER<br>(Include Area Code)<br>(805) 277-5440                                    | 22c. OFFICE SYMBOL<br>AFRPL/LKC |          |               |

DD FORM 1473, 83 APR

EDITION OF 1 JAN 73 IS OBSOLETE

SECURITY CLASSIFICATION OF THIS PAGE

the higher energy fraction for the charged pions makes hydrogen more attractive as the working fluid. It is therefore important to investigate possible means for efficient transfer of the pion energy to the working fluid.

# TABLE OF CONTENTS

| <u>SECTION</u>                                       | <u>PAGE</u> |
|--|-------------|
| 1. Introduction.....                                 | 1           |
| 1.1 Choice of Annihilating Materials.....            | 1           |
| 2. The Literature Search.....                        | 3           |
| 2.1 Procedures and Results.....                      | 4           |
| 2.2 Antiproton - Heavy Nucleus Research.....         | 4           |
| 3. Annihilation Physics.....                         | 5           |
| 3.1 In Heavy Nuclei .....                            | 8           |
| 3.2 Nuclear Fragment Energies.....                   | 8           |
| 4. Theoretical Results for Uranium.....              | 10          |
| 4.1 Fraction of Energy to Charged Fragments.....     | 12          |
| 5. Experimental Results for Uranium and Silicon..... | 14          |
| 5.1 Annihilation Cross Sections.....                 | 16          |
| 5.2 Fraction of Energy to Charged Fragments.....     | 18          |
| 6. Discussion and Conclusions.....                   | 20          |
| References.....                                      | 22          |
| Bibliography.....                                    | 24          |



## LIST OF TABLES AND FIGURES

| <u>Tables</u>  | <u>Page</u> |
|--|-------------|
| Table 1. Characteristics of pions from $p + \bar{p}$ .....6<br>annihilation at rest.   | 6           |
| Table 2. Possible energies that may be available.....9<br>to nuclear fragments following annihilation<br>of an antiproton at rest in a heavy nucleus.  | 9           |
| Table 3. Theoretical characteristics of nuclear.....11<br>fragments, particles, and the residual<br>nucleus resulting from the encounter of a<br>175-MeV antiproton with a $^{238}\text{U}$ nucleus.           | 11          |
| Table 4. Theoretical kinetic energy of heavy, charged.....14<br>fragments (protons and heavier) resulting<br>from annihilation of an antiproton in a<br>uranium-238 nucleus.                                   | 14          |
| Table 5. Experimental results for the kinetic energy.....19<br>of heavy charged fragments (p, d, t, $^3\text{He}$ , $\alpha$ )<br>from the annihilation of an antiproton in<br>silicon and uranium-238 nuclei. | 19          |

### Figures

|  |    |
|--|----|
| Figure 1. Kinetic energy spectrum of charged pions.....7<br>from antiproton-proton annihilation at rest.   | 7  |
| Figure 2. Momentum-differential cross sections for.....15<br>proton production in the annihilation of<br>180 MeV antiprotons in silicon and uranium-238. | 15 |
| Figure 3. Experimental antiproton absorption cross.....17<br>sections as functions of the antiproton<br>momentum.  | 17 |



## 1. INTRODUCTION

Antiprotons colliding with a heavy nucleus annihilate with a proton or neutron in the nucleus. The numbers and energies of the particles emitted from the nucleus following annihilation are significant to the use of antiproton annihilation as a source of energy for spacecraft propulsion. An important quantity to consider is the fraction of the annihilation energy that goes into the kinetic energy of heavy charged particles emitted from the nucleus, such as protons and deuterons. A large value of this energy fraction facilitates transferral of the annihilation energy to a working fluid or plasma that forms the rocket exhaust. It has been the purpose of the work reported here to determine the energy fraction from relevant literature and to assess the significance of its value to annihilation propulsion. This purpose is important to the question of what form of matter should be employed in annihilation with antiprotons to achieve maximum efficiency in converting annihilation energy into propulsion energy.

### 1.1 Choice of Annihilating Materials

Consideration of matter-antimatter annihilation as an energy source for space propulsion has been taking place over the last several years. For details of the research, the reader is referred to the journal articles and reports by Forward, Morgan, Vulpetti, and Massier listed in the bibliography. Comprehensive research results are found in Air Force Rocket Propulsion Laboratory document, AFRPL-TR-85-034 (Forward) and in Volume 35 (1982) of the Journal of the British Interplanetary Society.

Matter-antimatter annihilation produces the greatest amount of energy per unit mass of propellant of any known possible means of propulsion. The form of antimatter most often considered for annihilation consists of antiprotons, which are the antiparticles to ordinary protons. The antiprotons might be contained in solid antihydrogen, each atom (or antiatom) of which is composed of an antiproton and a positron (an antielectron). Antiprotons are the preferred form of antimatter because each antiproton annihilation produces about 2000 times as much energy per particle as positron annihilation with electrons, and antiprotons are much easier to produce than antinuclei

(composed of antiprotons and antineutrons). Antimatter of any other form involves antiparticles that are not stable; they would decay in storage before being used. In contrast, the form of matter to employ for annihilation with the antiprotons is less certain. Antiprotons annihilate with the protons and neutrons in the nuclei of matter atoms, but whether the nucleus should be a proton (nucleus of ordinary hydrogen), a heavy nucleus (e.g. of uranium), or something in between is not clear. The choice depends on how the annihilation energy manifests itself and on how it can be used to provide propulsion.

When protons or other light nuclei are employed, energetic pions (mesons with a mass of about  $1/7$  of the proton mass) are the principal product produced by the annihilation. Most of the pions are charged, and they can be directed by a magnetic field to produce thrust. However, their exhaust velocity is about 90% of the speed of light. Thus, for nearly all envisioned missions (not including interstellar flight) where the space craft velocity is much less than the speed of light, only a small fraction of the annihilation energy is transferred to the spacecraft; nearly all remains in the exhaust. To achieve a higher efficiency, it is therefore necessary to transfer the annihilation energy to a working fluid, of much higher mass than the antimatter, that has a much lower exhaust velocity but much greater thrust than the pions. The efficiency of energy transferral is affected by the form of matter chosen as the annihilation medium.

If the annihilation medium consists of an element of high atomic number, some of the pions produced by annihilation (of antiprotons with the protons and neutrons of the heavy nuclei of that element) will transfer their energy to the nucleus, and nuclear fragments (protons, neutrons, deuterons, and other light nuclei) will be emitted. The masses of the fragments are much greater than the pion mass, so it is much easier to transfer the energy of the charged fragments to a working fluid than to transfer the energy of the charged pions. (Transfer of energy of neutral particles of either kind is still more difficult.) Ease of transfer means less distance that a particle or fragment must travel through the working fluid to transfer its energy. Large transfer distances might require rocket motors (where the transferral would occur) that are much larger than those used for chemical propellants.

If the fraction of annihilation energy going into the kinetic energy of charged nuclear fragments for heavy nuclei is about the same or greater than the fraction of energy going into the kinetic energy of charged pions for light nuclei, then choosing a heavy element for the matter will be advantageous as opposed to hydrogen or another light element where essentially no fragments are produced. Therefore, it is important to determine the fraction of antiproton annihilation energy that goes into the kinetic energy of charged nuclear fragments when antiprotons annihilate in heavy nuclei. To that end, a literature survey was conducted to locate information from which the energy fraction could be determined.

## 2. THE LITERATURE SEARCH

The literature search was conducted to obtain information on the annihilation of antiprotons in nuclei with a mass equal to, or greater than, carbon. Occasionally the search area was broadened to include lighter nuclei. Of greatest interest was the annihilation of antiprotons in uranium-238 nuclei, the most massive nuclei of any naturally occurring isotope. It appeared likely that this annihilation would have the largest fraction of pion energy transferred to charged nuclear fragments, because the large size of the nucleus leads to a large distance of travel for the pions within the nucleus. Carbon was chosen as the "lower limit" on heavy nuclei for two reasons. First, on a logarithmic scale the carbon nuclear radius is about midway between the nuclear radii of uranium and a proton. Second, carbon is a fairly common target in antiproton experiments, so it appeared likely that some relevant experimental information would exist. Information was sought on annihilation in the hydrogen and helium isotopes because it is relevant to the basic antiproton-proton and antiproton-neutron annihilations.

## 2.1 Procedures and Results

The search was conducted by consulting reference documents, by using computer data bases, and by communicating with individuals in the field of high energy physics. The Physics Briefs, INSPEC, American Institute of Physics, NTIS, and DTIC data bases were searched for publications over the last five to eight years. This yielded about two hundred citations for examination. Approximately fifty of the cited publications were obtained, because they were relevant to determining the energy fraction. Many of these had references to relevant publications of twenty to thirty years ago which were also obtained. The relevant publications are included in the Bibliography.

The personal contacts produced much useful information, including the results of a calculation of antiproton - uranium-238 annihilation,<sup>(1)</sup> and experimental results for antiproton annihilation in uranium-238 and silicon.<sup>(2)</sup> These results, when combined with earlier work, allow determination of both theoretical and experimental values for the energy fraction.

## 2.2 Antiproton - Heavy Nucleus Research

Experimental research on antiproton - heavy nucleus annihilation is concentrated in two time periods. The first period began shortly after the discovery of the antiproton in 1955<sup>(3)</sup> and continued into the early 1960's. The second period began about four years ago and includes the present. That period followed a time of heightened theoretical interest in the consequences of the large energy deposited in nuclear matter by annihilation.<sup>(4)</sup> This energy can lead to interesting physical circumstances (described in terms of quarks and gluons) that can improve the understanding of strong forces. The initial experimental work involved the study of charged particle tracks in photographic emulsions<sup>(5)</sup> and bubble chambers<sup>(6)</sup> exposed to antiproton beams at Lawrence Berkeley National Laboratory and subsequently at Brookhaven National Laboratory. At the present time most experimental work is being conducted at the LEAR (Low Energy Antiproton Ring) facility in the Antiproton Complex at CERN (Centre Européenne pour la Recherche Nucléaire, now named

European Organization for Nuclear Research) near Geneva, Switzerland (e.g., Ref. 2).

Some values for the energy fraction were determined during the first period of experimental work<sup>(5,6)</sup> However, these values are based on observation of a small number of annihilations (compared to modern capabilities) and either do not apply to specific nuclei or do not include heavy charged particles and fragments of all significant energies and types. The literature search yielded no published values of the energy fraction from the second period of experimental work, but many of the results from recent experiments at LEAR on antiproton annihilation in heavy nuclei are not fully published.<sup>(7)</sup> Hence the possibility of modern published values of energy fractions in the near future. The various pieces of experimental and theoretical information that currently exist, however, may be used to calculate the fraction of annihilation energy transferred to the kinetic energy of heavy, charged particles (or nuclear fragments) in antiproton annihilation with uraniums-238 and silicon nuclei.

### 3. ANNIHILATION PHYSICS

When an antiproton ( $\bar{p}$ , single negative charge) annihilates with a proton ( $p$ , single positive charge) or neutron ( $n$ , neutral) in a heavy nucleus, a number of lighter particles are produced. The number and types of particles vary from annihilation to annihilation. In nearly all annihilations, however, the particles are types of mesons, with pi-mesons (pions) being most likely. Most of the other mesons are "resonances," so called because they are short-lived excited states of more stable mesons, which decay into other particles (pions, to a large degree, in roughly  $10^{-23}$  to  $10^{-19}$  s) by way of strong forces. A small number of K-mesons (kaons) are produced directly or by decays of the short-lived mesons. The pions and kaons decay via weak and electromagnetic forces so they have much longer lifetimes (about  $10^{-16}$  to  $10^{-8}$  sec). The kaons constitute only about 2% of the decay products.

The three above "forces" plus gravity are the four fundamental forces of nature. In decreasing order of strength: strong forces are the nuclear

forces that bind protons and neutrons in nuclei and give rise to fission and fusion energy; electromagnetic forces bind electrons to nuclei and are responsible for most physical, chemical, and biological phenomena; weak forces cause the radioactive decay of nuclei and the decay of some fundamental particles; and gravitational forces are responsible for weight and the motion of celestial bodies.

For many purposes including antiproton-annihilation propulsion, the antiproton annihilation products may be taken to be entirely pions, including charged pions ( $\pi^+$  and  $\pi^-$ ) and neutral pions ( $\pi^0$ ). Characteristics of the pions from the annihilation of an antiproton at rest with a proton (i.e., when the relative velocity is nearly zero) are given in Table 1. The energy spectrum of the charged pions is shown in Fig. 1.

The pions eventually decay. The  $\pi^+$  and  $\pi^-$  decay via weak forces into muons and neutrinos (the muons then decay into electrons, positrons, and neutrinos) while the  $\pi^0$ 's decay via electromagnetic forces into photons (gamma rays).

The cross section for the annihilation of an antiproton with a proton is about 1.3 times the cross section for annihilation with a neutron for equal, low-incident energies.<sup>(8)</sup> This ratio was determined from observing annihilations of protons and neutrons in nuclei, but it also applies to free protons and neutrons. In annihilation with a neutron, the number of negative

Table 1. Characteristics of pions from  $p + \bar{p}$  annihilation at rest (information derived from Ref. 12 and other sources in the Bibliography). Total annihilation energy = mass energy of  $p + \bar{p}$  = 1876.61 MeV. About 4% of the annihilation energy goes into other particles (mainly kaons).

| particle | mean number of particles of each type per annihilation | mean kinetic energy per particle [MeV] | mass energy of particle [MeV] | mean life-time [s]    |
|----------|--|--|-------------------------------|-----------------------|
| $\pi^+$  | 1.60   | 235                                    | 139.58                        | $2.60 \times 10^{-8}$ |
| $\pi^-$  | 1.6  | 202                                    | 134.98                        | $9.1 \times 10^{-17}$ |
| $\pi^0$  | 1.4  | 225                                    | 139.58                        | $2.6 \times 10^{-8}$  |

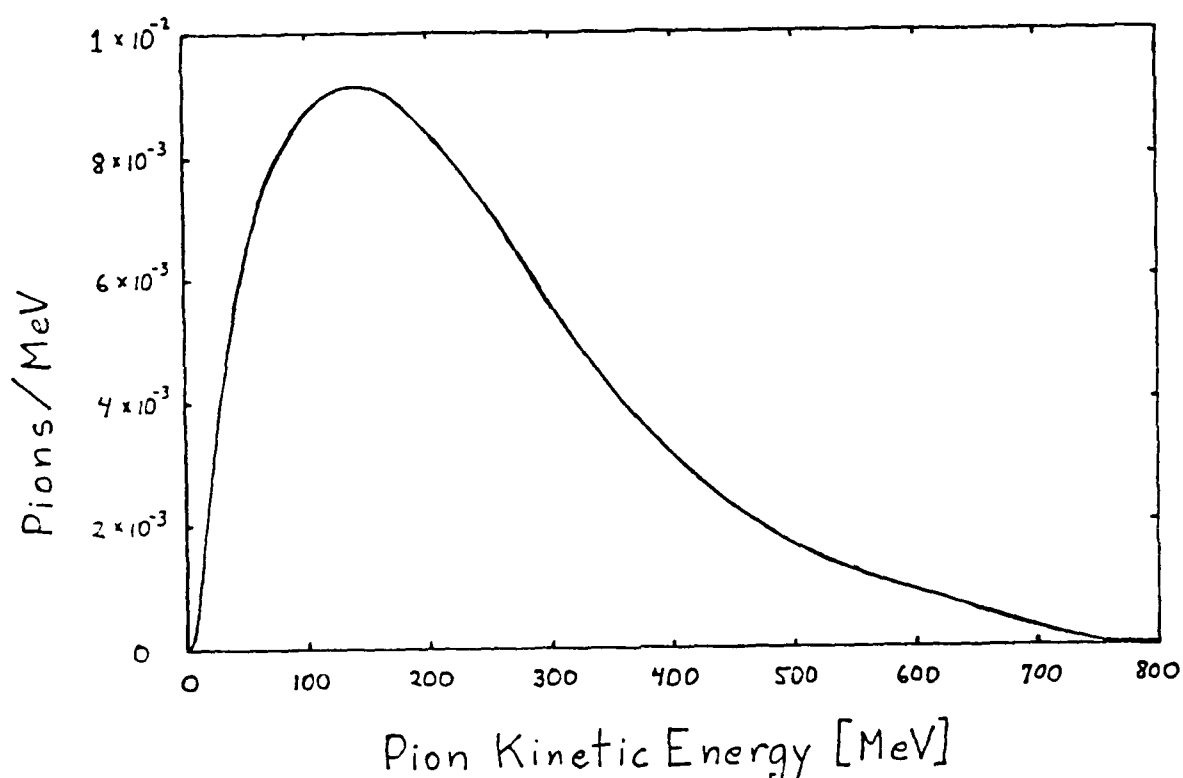


Fig. 1. Kinetic energy spectrum of charged pions from antiproton-proton annihilation at rest. Adapted from Ref. 17 for 0-600 MeV, from Ref. 18 for 600-700 MeV, and extrapolated to calculated cut-off for annihilation into  $\pi^+ + \pi^- + \pi$  for greater than 700 MeV. The spectrum is normalized to a single annihilation, thus the area under the curve is 3.00, the mean number of charged pions per annihilation.

pions produced is always one more than the number of positive pions since charge must be conserved. The mean relative proportion of charged pions is somewhat greater than for annihilation with a proton, but the mean energies of the pions are about the same.

### 3.1 In Heavy Nuclei

When an antiproton with an energy of a few hundred MeV or less strikes a heavy nucleus, the probability for annihilation with either a proton or neutron is so high that annihilation occurs on the surface of the nucleus. The annihilation pions are emitted isotropically in the center of mass system of the annihilating particles, so about one-half of the pions will enter large nuclei. The cross section (or probability) for interaction of these pions with the nucleons (neutrons and protons) is sufficiently large that they will interact with one or more nucleons before exiting the nucleus, if they indeed get out. A pion may scatter from a nucleon or may be absorbed by a nucleon. When absorption occurs, the excited nucleon will usually re-emit a pion. In both cases, some of the kinetic energy of the pion is transferred to the nucleon. The energy transferred is usually well above the 5 - 10 MeV mean binding energy with which each nucleon is bound within the nucleus. Thus, a number of nuclear fragments, which are individual nucleons or combinations of nucleons, will be emitted from the nucleus. In addition, the nucleus will be left in an excited state. If the residual nucleus is sufficiently heavy to be unstable to fission, then it is likely that fission will follow the cascade of the pions through the nucleus.<sup>(9)</sup> The fission energy transferred to the kinetic energy of the two roughly equal nuclei that result is about 175 MeV.<sup>(10)</sup>

### 3.2 Nuclear Fragment Energies

Most of the fragments emitted will be individual protons and neutrons.<sup>5,6</sup> A few will be deuterons (designated d, one p and one n), tritons (designated t, one p and two n's), helium-3 nuclei (designated  $^3\text{He}$ , two p's and one n), and alphas (designated  $\alpha$ , two p's and two n's). The total kinetic energy transferred to these fragments depends on the dynamics of the pion-nucleon interactions and on the size of the nucleus. Some possible energies that might be available to the fragments are given in the first two columns of



Table 2. The first column considers the kinetic energy of one-half of the pions, and in the second column the mass energy of these pions is also included. The last row of both columns additionally includes the kinetic energy of the fission fragments that may be present from fission of the residual nucleus. Table 2 also gives the kinetic energy of the charged pions produced by annihilation with a proton (about the same as for annihilation in other light nuclei consisting of only a few nucleons).

It is apparent from the figures in Table 2 that if a high fraction of the available energy is transferred to charged nuclear fragments when annihilation occurs in a heavy nucleus, then the kinetic energy of these fragments is comparable to the kinetic energy of the charged pions from annihilation with a proton (hydrogen nucleus). It is this possible fact that makes heavy elements potentially attractive as an annihilation medium and working fluid for antiproton annihilation (and led to this study), since transferral of energy to the working fluid with the relatively heavy charged nuclear fragments is easier to accomplish than with the relatively lighter pions.<sup>(11)</sup> It will be seen in the following two sections, however, that only a fairly small fraction of the available energy goes into the kinetic energy of charged fragments. Much goes into reemitted pions and a significant amount into the kinetic energy of neutrons.

Table 2. Possible energies that may be available to nuclear fragments following annihilation of an antiproton at rest in a heavy nucleus compared to the kinetic energy of charged pions produced by proton-antiproton annihilation. Small contribution from kaons not included.

|  | <u>Annihilation in heavy nucleus</u> |                              | <u>Annihilation with a proton</u>      |
|--|--------------------------------------|------------------------------|--|
|  | <u>Half of pion kinetic energy</u>   | <u>Same with mass energy</u> | <u>Kinetic energy of charged pions</u> |
| Energy [MeV]   | 556                                  | 900                          | 705                                    |
| Fraction of annihilation energy                            | 0.30                                 | 0.48                         | 0.38                                   |
| Fraction with kinetic energy of fission fragments included | 0.39                                 | 0.57                         | -----                                  |

#### 4. THEORETICAL RESULTS FOR URANIUM

Michael R. Clover of Los Alamos National Laboratory provided the output of a computer run that gives the details in the Intra-nuclear Cascade Model of  $\bar{p}$  scattering by a uranium-238 nucleus.<sup>(1)</sup> These results are similar to, but more detailed than, those reported by Clover et al. for the same problem.<sup>(12)</sup>

The code performed a monte carlo simulation of encounters between a 175 MeV (lab energy) antiproton and a  $^{238}\text{U}$  nucleus (at rest in the Lab frame). The number of encounters was 7029 of which 4950 were inelastic (the  $\bar{p}$  changes form or loses energy in the center-of-mass frame). The  $\bar{p}$  did not pass close enough to the nucleus for inelastic processes to occur in the other encounters. Nearly all of the inelastic encounters involved  $\bar{p}$  annihilation, however in 244 inelastic cases the  $\bar{p}$  remained intact but lost energy to individual nucleons, and in 17 cases charge exchange occurred. In these latter cases the antiproton struck a proton and an anti-neutron plus a neutron were produced, with the anti-neutron then leaving the nucleus without annihilating. In charge exchange, the antiproton gives its negative charge to a proton and becomes an antineutron ( $\bar{n}$ ), while the negative charge neutralizes the positive charge of the proton making it a neutron.

The code output contains the number and energy spectra of the nuclear fragments, pions, kaons, and other particles produced in the inelastic collisions. From the spectra, the mean energy for each particle or fragment may be calculated. The numbers and mean energies are give in Table 3. A correction was applied to these quantities to obtain similar quantities for the case of interest here in which only annihilations occur. Table 3 also contains those results, along with the mean total kinetic energy given to each particle or fragment type for a single annihilation. The particle and fragment numbers allow calculation of the mass energy of the created particles (pions and kaons) and the nuclear binding energy lost (potential energy gained) when the nucleons and fragments leave the nucleus. These are also in Table 3. The code gives information on the distribution of nuclear states following annihilation from which the mean excitation energy of the residual nucleus was determined. That value is shown in Table 3 where it is added with

Table 3. Theoretical characteristics of nuclear fragments, particles, and the residual nucleus resulting from the encounter of a 175-MeV antiproton with a  $^{238}\text{U}$  nucleus based on Ref. 1. Statistical errors are not shown.

| particle or fragment | Inelastic Scattering  |  | Annihilations only           |  |   |
|----------------------|-----------------------|--|------------------------------|--|---|
|                      | mean number per event | mean kinetic energy per particle or fragment [MeV] | mean number per annihilation | mean kinetic energy per particle or fragment [MeV] | total kinetic energy for each type per annihilation [MeV] |
| p                    | 1.419                 | 101.2  | 1.498                        | 101.7  | 152.4   |
| n                    | 3.608                 | 72.4   | 3.809                        | 72.8   | 277.4   |
| $\pi^+$              | 0.758                 | 212.   | 0.801                        | 213.   | 170.6   |
| $\pi^0$              | 1.137                 | 197.   | 1.201                        | 198.   | 238.  |
| $\pi^-$              | 1.206                 | 193.   | 1.273                        | 194.   | 247.  |
| $K^+$                | 0.030                 | 84.  | 0.032                        | 85.  | 2.7   |
| $K^0$                | 0.068                 | 65.  | 0.071                        | 66.  | 4.7   |
| $K^-$                | 0.039                 | 99.  | 0.041                        | 99.  | 4.1   |
| $\bar{p}$            | 0.049                 | 146.   | -                            | -  | -   |
| $\bar{n}$            | 0.003                 | 134.   | -                            | -  | -   |
| d                    | 0.347                 | 77.9   | 0.366                        | 78.4   | 28.7  |
| t                    | 0.266                 | 62.3   | 0.280                        | 62.6   | 17.6  |
| $^3\text{He}$        | 0.163                 | 65.6   | 0.172                        | 66.0   | 11.4  |
| $\alpha$             | 0.044                 | 54.3   | 0.047                        | 54.6   | 2.6   |

|  |           |
|--|-----------|
| Total kinetic energy to particles/fragments          | 1156. MeV |
| Total mass energy of $\pi$ 's & $K$ 's               | 523.      |
| Nuclear recoil kinetic energy                        | 1.        |
| Nuclear binding energy lost (approx.)                | 59.       |
| Mean excitation energy of residual nucleus (approx.) | 296.      |
| Total Discrepancy                                    | 17.       |
| p + $\bar{p}$ mass energy + $\bar{p}$ kinetic energy | 2052. MeV |

the other energies in an attempt to reproduce the total input energy, the sum of the annihilation energy and the incident kinetic energy of the antiproton. This sum falls 17 MeV short of the input energy. One possible explanation involves the assumption employed that the mean binding energy per nucleon is the same in the possible residual nuclei as in  $^{238}\text{U}$ , whereas it is likely somewhat less. Another possible explanation is a small systematic

error in the numerical integration of the energy spectra to obtain the mean energies of the particles and fragments. In any case, the discrepancy is less than one percent of the total energy, and hence negligible for our purposes.

The residual nucleus, after annihilation, is in a highly excited state from which more fragments (mainly nucleons) will be emitted by evaporation, along with the emission of gamma rays. The contribution of evaporation to the total energy in emitted fragments may be estimated by considering Fig. 4 of Ref. 13. Here it is shown that the energy deposited by evaporated protons following antiproton annihilation in a 1 mm thick slab of silicon is about the same as the energy deposited by the protons emitted during the initial intra-nuclear cascade. Using this fact and a formula for the rate of energy deposition for charged particles in matter<sup>(14)</sup>, it can be determined that the total kinetic energy in the evaporated protons is less than one-tenth that of the initial cascade protons when the initial energy of the evaporated protons is 25 MeV. If that initial energy is 10 MeV, then the ratio of the energies is one-hundredth. The reason for the nearly equal deposition of energy is that the lower-energy, evaporated protons deposit their energy much more rapidly than the faster cascade protons. Since the mean energy of each evaporated proton is probably around 10 MeV and certainly less than 25 MeV,<sup>(15)</sup> the contribution of evaporated protons is insignificant. The same is very likely true for other fragments, so evaporated fragments may be neglected.

#### 4.1 Fraction of Energy to Charged Fragments

The results of Table 3 give 213 MeV for the total kinetic energy of heavy charged fragments (p,d,t,<sup>3</sup>He, $\alpha$ ) following annihilation of a 175 MeV antiproton in a <sup>238</sup>U nucleus. In an antiproton annihilation rocket engine, the antiprotons will most likely be at much lower energy. They may be stored at temperatures near absolute zero<sup>(11)</sup> ( $10^{-5}$  to  $10^{-4}$  eV, perhaps) and extracted and transported to the rocket engine at energies well under 1 MeV. As far as the intra-nuclear dynamics is concerned, the annihilation of such an antiproton amounts to annihilation "at rest". More energy is transferred from the pions to the fragments for a  $\bar{p}$  energy of 175 MeV than a  $\bar{p}$  energy of nearly zero. This is because center-of-mass motion at 175 MeV tilts the pion

distribution forward resulting in more pions entering the nucleus and because the pions share the additional 175 MeV of energy.<sup>(5,6)</sup> Available experimental information allows determination of the ratio of kinetic energy transferred to charged fragments for annihilation at rest to the same quantity for annihilation at 175 MeV.

From Ref. 6 the value of the above ratio, when comparing annihilations at 25 MeV and 120 MeV is  $0.84 \pm 0.08$ . Assuming an inverse linear dependence on energy, this gives  $0.73 \pm .12$  for the ratio from 0 MeV to 175 MeV. From Ref. 5 the ratio is  $0.57 \pm 0.06$  for zero to 166 MeV. When similarly extrapolated to the case of 0 MeV to 175 MeV, this latter ratio becomes  $0.56 \pm 0.06$ . Combining the two ratios for 0 MeV to 175 MeV with somewhat more weight for the latter, the value of  $0.62 \pm 0.10$  is obtained as a single value for the ratio, where the large error reflects the discrepancy in the two ratios. Application of this single ratio yields  $130 \pm 20$  MeV as the theoretical value for the kinetic energy of the heavy charged fragments (p,d,t,<sup>3</sup>He, $\alpha$ ) resulting from antiproton annihilation at rest in a <sup>238</sup>U nucleus. That value, along with related quantities is given in Table 4. It may be seen from Table 4 that annihilation of antiprotons in heavy nuclei is less attractive than indicated by earlier estimates of an energy fraction of about 0.5.<sup>(11)</sup> A similar conclusion occurs when experimental information on  $\bar{p}$ -heavy nucleus annihilation is considered in the following section. These results apply whether or not the <sup>238</sup>U nucleus or proton are free or in uranium and hydrogen atoms (or molecules) since they pertain to events per annihilation. The annihilation rate for antiprotons at energies around a few eV or less is much higher when the annihilating medium consists of atoms or molecules than free (bare) nuclei<sup>(11)</sup> (see also entries in Bibliography by Morgan and Hughes).

Table 4. Theoretical kinetic energy of heavy, charged fragments (protons and heavier) resulting from annihilation of an antiproton in a uranium-238 nucleus (based on Ref. 1) compared to the kinetic energy of charged pions in  $p + \bar{p}$  annihilation at rest. Quantities for annihilation at rest are obtained by applying a factor to those at 175 MeV. The factor is obtained from experimentally based information in Refs. 5 and 6.

| annihilating nucleus                                       | $^{238}\text{U}$      | $^{238}\text{U}$               | p                 |
|--|-----------------------|--------------------------------|-------------------|
| $\bar{p}$ energy [MeV]                                     | 175                   | 0<br>(at rest)                 | 0<br>(at rest)    |
| energy to charged fragments/pions [MeV]                    | 210<br>(to fragments) | $130 \pm 20$<br>(to fragments) | 705<br>(to pions) |
| fraction of annihilation (plus any incident) energy        | 0.10                  | $0.07 \pm .01$                 | 0.38              |
| fraction with kinetic energy of fission fragments included | 0.19                  | $0.16 \pm .01$                 | -----             |

## 5. EXPERIMENTAL RESULTS FOR URANIUM AND SILICON

Among the results of LEAR experiment PS187 reported in Ref. 2 are momentum-differential cross sections for proton production in 180 MeV antiproton annihilation with Si and  $^{238}\text{U}$  nuclei (see Fig. 3 in Ref. 2, reproduced here as Fig. 2). By reading values from the graph, the cross sections may be integrated to obtain the total cross section for proton production as well as the mean energies of the protons. That was done under the assumption that there is no significant contribution to the total kinetic energy of the protons for momenta below 250 MeV/c (kinetic energy = 33 MeV) which is the low momentum cutoff on the graph. For  $^{238}\text{U}$ , the cross section is 5800 mb (millibarn =  $10^{-27} \text{ cm}^2$ ) and the mean kinetic energy is 100 MeV. For Si the cross section is 1350 mb and the mean kinetic energy is 110 MeV. To obtain the kinetic energy of the protons for each annihilation requires knowing the number of protons emitted per annihilation. That number is equal to the proton production cross section divided by the annihilation cross section, but the annihilation cross sections are not given in Ref. 2 (nor in Ref. 13 on the same experiment). They can, however, be obtained from other experimental information.

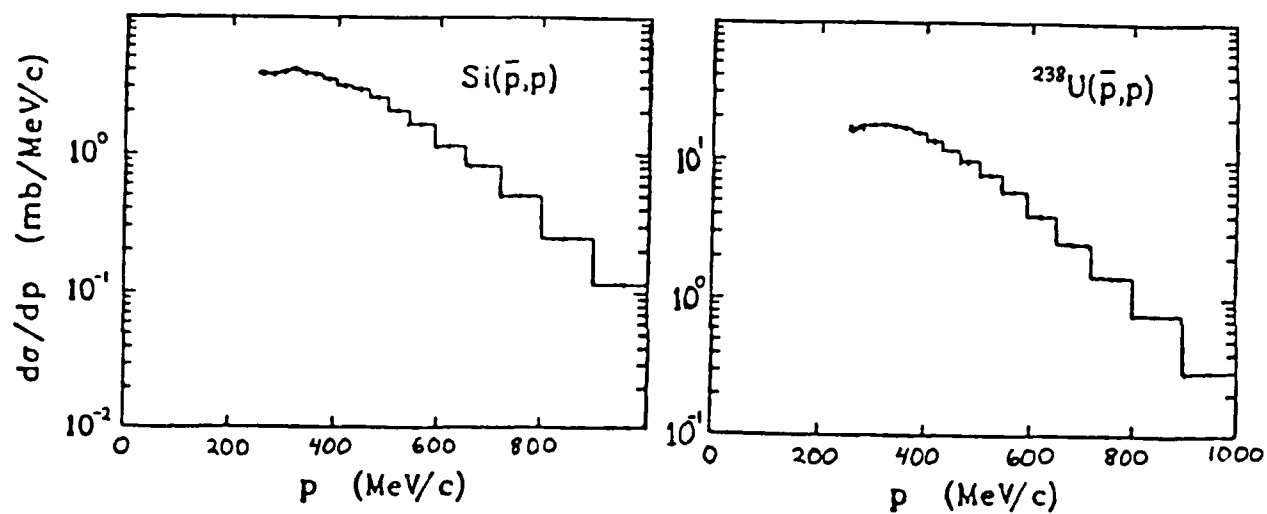


Fig. 2. Momentum-differential cross sections for proton production in the annihilation of 180 MeV antiprotons in silicon and uranium-238. Note, the symbol  $p$  is used both to designate a proton and for the momentum of a particle (here, a proton). Taken from Fig. 3 of Ref. 2.

## 5.1 Annihilation Cross Sections

Nakamura et al.<sup>(16)</sup> give measured values of  $\bar{p}$  annihilation cross sections in carbon, aluminum, and copper for antiproton energies between 100 MeV and 350 MeV. At such energies, the cross sections are the same for free nuclei and ones in atoms. They call the cross sections "absorption cross sections", but it is evident that they are essentially equal to annihilation cross sections because the only other inelastic process included was charge exchange, which accounts for only about 0.3% of the absorption cross section. Besides their own values, Nakamura et al. also give annihilation cross sections in their Fig. 2 from other measurements for higher energies. That figure is reproduced here as Fig. 3. Because of the fairly simple dependence of annihilation cross section on atomic number of the nucleus and on the momentum or energy of the antiproton, the measurements of Nakamura et al., along with the other measurements in Fig. 3 can be used to determine the antiproton annihilation cross sections for silicon and uranium-238.

For antiproton energies high enough that the DeBroglie wave length of the antiproton is small compared to the nuclear radius (true for the momenta in Fig. 3), the annihilation cross section is approximately proportional to the geometric cross section of the nucleus. In fact, this proportionality is a near-equality for antiproton energies of a few hundred MeV (momenta of several hundred MeV/c) since the proton has a near-unity probability of annihilating once it encounters the nucleus. At higher energies, the nucleus begins to become transparent to the antiproton. At much lower energies, the wavelength of the antiproton becomes large and the quantum-mechanical spreading of the antiproton destroys the simple geometric picture.

In this simple but fairly accurate view, the annihilation cross section has the form

$$\sigma_a = \pi C A^{2/3} p^{-8} \quad (1)$$



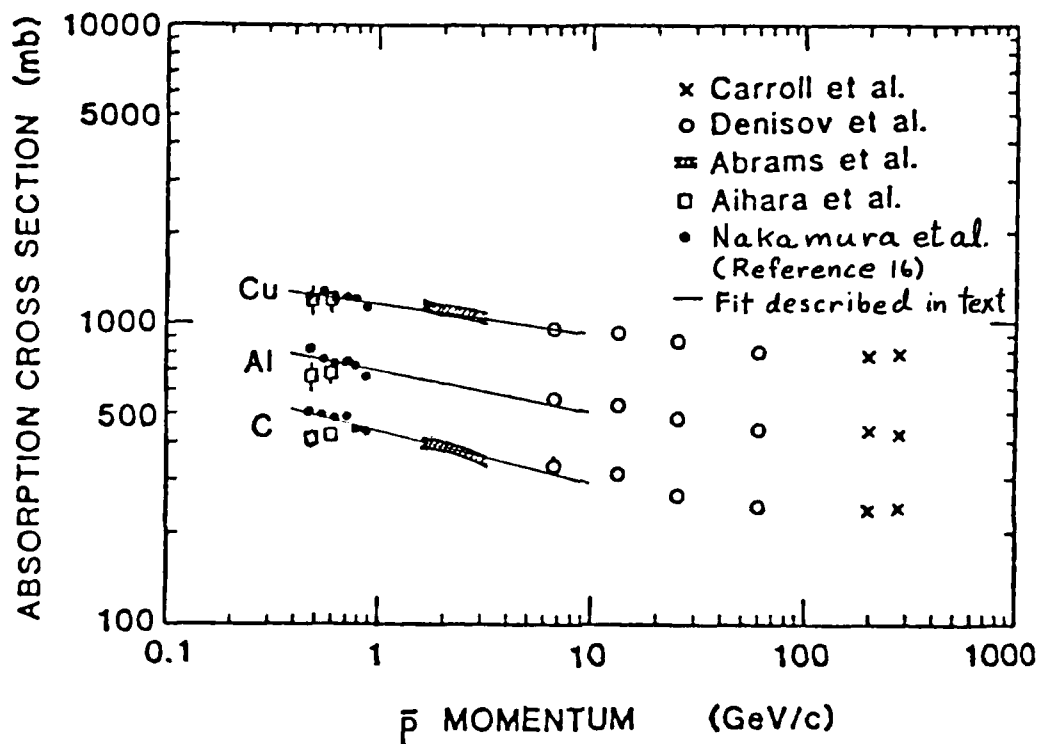


Fig. 3. Experimental antiproton absorption cross sections as functions of the antiproton momentum, from Fig. 2 of Ref. 16. The absorption cross section is essentially equal to the annihilation cross section. The solid curves are the fit described in the text. See Ref. 16 for full reference information on the experimental work.

where  $A$  is the atomic number of the nucleus (the area of the nucleus is proportional to  $A^{2/3}$ ),  $p$  is the momentum of the antiproton (same symbol as used to designate a proton), and  $C$  and  $\beta$  are constants. The factor  $p^{-\beta}$  can account, at least to a degree, for both the transparency of the nucleus as  $p \rightarrow \infty$  and the spreading of the antiproton as  $p \rightarrow 0$ . A more accurate form for  $\sigma_a$  may be obtained by allowing  $\beta$  to depend somewhat on  $A$ , as appears must be the case from Fig. 3, and by replacing  $C A^{2/3}$  with  $(C^{1/2} A^{1/3} + d)^2$ . The constant  $d$  accounts for the fact that annihilation may still take place out to some fixed distance beyond the mean radius of the nucleus. When such an altered form is used to fit the data of Fig. 3 for momenta less than 10 GeV/c, the results is

$$\sigma_a = \pi(1.35 A^{1/3} + 0.83)^2 \left( \frac{p}{600 \text{ MeV/c}} \right)^{-0.5A^{-0.4}} \times 10^{-26} \text{ cm}^2. \quad (2)$$

Equation 2 gives the solid lines shown in Fig. 3 and gives a value of 2650 mb for  $^{238}\text{U}$  with an antiproton energy of 175 MeV. This value is within 5% of the theoretical value of 2500 mb.<sup>(1)</sup> Equation 2 should not be employed for nuclei lighter than carbon. (For  $\bar{p} + p$  it gives about twice the known value.)

## 5.2 Fraction of Energy to Charged Fragments

Equation 2 may now be employed to obtain the annihilation cross sections for a 180 MeV antiproton (momentum = 608 MeV/c) on Si and  $^{238}\text{U}$  (as in LEAR experiment PS187) and thence the numbers and energies of the protons produced by the annihilations. For  $^{238}\text{U}$ ,  $\sigma_a = 2650$  mb so the number of protons per annihilation is 2.19 and their kinetic energy per annihilation is 219 MeV. To this must be added the energy of the other heavy charged particles which are not considered in Ref. 12 or 13. However, using the ratio, 0.39, of the combined kinetic energy of d, t,  $^3\text{He}$ , and  $\alpha$  to the kinetic energy of the protons from the theoretical calculation of the previous section (since no experimental value is known), the total kinetic energy of the heavy charged fragments is found to be 305 MeV. If the procedure in the previous section is used to extrapolate the energy to annihilations at rest, the factor to be

employed is  $0.61 \pm 0.10$ , and the result for the kinetic energy of the charged fragments per annihilation is  $190 \pm 30$  MeV. For silicon,  $\sigma_a = 760$  mb, the mean number of protons is 1.75, their kinetic energy per annihilation is 195 MeV, and the kinetic energy of the charged fragments is 270 MeV (for a  $\bar{p}$  energy of 180 MeV). When extrapolated to annihilation at rest, the kinetic energy of charged fragments is  $165 \pm 25$  MeV per annihilation. These results are summarized in Table 5.

The experimental values for fraction of annihilation energy in  $^{238}\text{U}$  going into heavy charged fragments shown in Table 5 are somewhat higher than the theoretical results of Table 4, but the results are, nevertheless, in fair

Table 5. Experimental results for the kinetic energy of heavy charged fragments (p,d,t, $^3\text{He}$ , $\alpha$ ) from the annihilation of an antiproton in silicon and uranium-238 nuclei compared to the kinetic energy of the charged pions from  $\bar{p} + p$  annihilation (energies in MeV). The relative contribution for charged fragments other than protons is based on the theoretical results of Ref. 1.

|                  |   | $\bar{p}$ incident energy |  |                    |                  |
|------------------|---|---------------------------|--|--------------------|------------------|
| nucleus          |   | 180 MeV                   |  | at rest            |                  |
|                  |   | w/o fission energy        |  | w/o fission energy | w/fission energy |
| Si               | kinetic energy of charged fragments                 | 270                       |  | $165 \pm 25$       | -----            |
|                  | fraction of annihilation (plus any incident) energy | 0.13                      |  | $0.09 \pm 0.015$   | -----            |
| $^{238}\text{U}$ | kinetic energy of charged fragments                 | 305                       |  | $190 \pm 30$       | $365 \pm 30$     |
|                  | fraction of annihilation (plus any incident) energy | 0.15                      |  | $0.10 \pm 0.015$   | $0.19 \pm 0.015$ |
| p                | kinetic energy of charged pions                     | -----                     |  | 705                | -----            |
|                  | fraction of annihilation energy                     | -----                     |  | 0.38               | -----            |

agreement. The higher values are still not as high as previously supposed in Ref. 11. Thus, annihilation of an antiproton in a heavy nucleus is not as attractive as suggested in that report.

The values in Table 5 show only a small relative difference between the kinetic energies of charged fragments from silicon and uranium-238. Therefore, if it were desirable to employ antiproton-heavy nucleus annihilation, medium-weight nuclei would work about as well as the heaviest nuclei when fission energy is not considered.

## 6. DISCUSSION AND CONCLUSIONS

Previous studies have indicated a possible difficulty in coupling the antiproton annihilation energy to a working fluid that forms the rocket exhaust in antiproton annihilation propulsion.<sup>(11)</sup> This difficulty is due to the large distance that the relatively light charged pions from antiproton-proton annihilation must travel to transfer their energy, compared to possible distances that can be travelled by the pions (or decay muons) while confined in a magnetically contained plasma of a few meters size. If annihilation were to occur in a heavy nucleus of a working fluid of high atomic number rather than in hydrogen, part of the annihilation energy would go into the kinetic energy of charged, heavy nuclear fragments emitted from the nucleus such as protons, deuterons, etc., whose energy could be more readily transferred to the working fluid. The purpose of this study was to determine the fraction of annihilation energy that goes into the kinetic energy of the charged, heavy nuclear fragments.

A literature survey yielded sufficient theoretical and experimental information from which the energy fraction could be determined. Its value is about 10% for nuclei as massive as silicon or greater. It is the same or less for lighter nuclei. If fission energy is included for very heavy nuclei, the value of the energy fraction is about 20%. Both of these figures are less than an earlier estimate of about 50%,<sup>(11)</sup> and they are significantly less than the fraction of the annihilation energy, 38%, that goes into the kinetic energy of charged pions in antiproton-proton annihilation. If the kinetic

energy of the charged pions is included, the corresponding energy fractions are about 30% for nuclei as heavy as silicon or greater and 40% for very heavy nuclei when fission is considered. These figures are still less than, or not significantly greater than, 38%. Their low values are mainly a consequence of the fact that a significant amount of the annihilation energy goes into the kinetic energy of emitted neutrons. I assume that the amount of energy that can transfer from neutral particles to the working fluid is insignificant.

Thus, for plasma combustion chambers of a few meters size, annihilation of antiprotons in heavy nuclei does not offer an advantage over annihilation with protons if an effective means can be found to couple the kinetic energy of the charged pions from annihilation with protons to a working fluid. If such a means cannot be found, then annihilation in heavy nuclei will allow transferral of up to 10-20% of the annihilation energy to the working fluid.

Since annihilation with protons has the potential for giving two to four times as much energy to the working fluid as annihilation in heavy nuclei, it is important to investigate the transferral of energy from charged pions to a working fluid in more detail. One must determine the mechanisms that would allow efficient transferral of energy in a combustion chamber with dimensions of a few meters, inside a magnetic field (to contain the charged pions), and at attainable plasma densities. In the work for Ref. 11, estimates of energy transferral from the charged pions to nuclei and bound electrons were employed. These estimates must be replaced by accurate calculations, and transferral of energy to free electrons in the plasma (electron drag) must be added. In addition, it is important to determine the mean distance traveled by charged pions (or their decay muons) before they exit the chamber, as a function of the plasma density and the strength and configuration of the magnetic field.

# REFERENCES

1. Michael R. Clover, Los Alamos National Laboratory (LANL), private communication consisting of a computer run of a modified version of the ISABEL intra-nuclear cascade code of Yariv and Fraenkel. See Y. Yariv and Z. Fraenkel, Phys. Rev. C24, 488 (1981) and C20, 2227 (1979) as well as Ref. 12. The computer run was carried out on LANL worker machine V at 10:42:52 on November 5, 1985.
2. P.L. McGaughey, K.D. Bol, M.R. Clover, R.M. DeVries, N.J. DiGiacomo, J.S. Kapustinsky, G.R. Smith, J.W. Sunier, W.E. Sondheim, Y. Yariv, M. Buenerd, J. Chauvin, D. Lebrun, P. Martin, and J.C. Dousse, Antiproton Annihilations in Nuclei, in "Hadronic Probes and Nuclear Interactions," AIP Conference Proceedings, No. 133, Ed. J.R. Comfort, W.R. Gibbs, B.G. Ritchie, HIPANI Conference, Tempe, Arizona, March 10-14, 1985, American Institute of Physics, New York, 1985, p.41.
3. O. Chamberlain, E. Segrè, C. Wiegand, and T. Ypsilantis, Phys. Rev. 100, 947 (1955).
4. See, for instance, the articles beginning on Page 495 in Section II (or referred to there) of Physics at LEAR with Low-Energy Cooled Antiprotons, Ed. U. Gastaldi and R. Klapisch, Plenum Press, New York, 1984.
5. A.G. Ekspong, A. Frisk, S. Nilsson, and B.E. Ronne, Antiproton Annihilation in Complex Nuclei, Nuclear Physics 22, 353 (1961).
6. L.E. Agnew, Jr., T. Elioff, W.B. Fowler, R.L. Lander, W.M. Powell, E. Segrè, H.M. Steiner, H.S. White, C. Wiegand, and T. Ypsilantis, Antiproton Interactions in Hydrogen and Carbon below 200 MeV, Phys. Rev. 118, 1371 (1960).
7. As evidenced, for example, by abstracts included in the portion of Sect. C following C29, abstracts included in Sect. G, and Abstract N10 of PANIC, Particles and Nuclei - Tenth International Conference, Book of Abstracts, Vol. 1 and 2, Heidelberg, July 30 - August 3, 1984, E.F. Güttner, B. Povh, and G. zu Putlitz, editors.
8. W. Chinowsky and G. Kojian, Antiproton Annihilation at Rest in Deuterium, Il Nuovo Cimento 43A, 684 (1966); V. Barnes et al., in Proceedings of the Twelfth International Conference on High Energy Physics, Dubna, USSR, 1964 (Atomizdat., Moscow, 1966), Vol. I, p.731.
9. S. Polikanov, letter to author, November 28, 1985.
10. American Institute of Physics Handbook, third edition, Ed. D.E. Gray, McGraw Hill, New York, 1972, p. 8-268.
11. D.L. Morgan, Jr., Concepts for the Design of an Antimatter Annihilation Rocket, Journal of the British Interplanetary Society 35, 405 (1982).
12. M.R. Clover, R.M. DeVries, N.J. DiGiacomo, and Y. Yariv, Low Energy Antiproton-Nucleus Interactions, Phys. Rev. C26, 2138 (1982)

13. J.W. Sunier, K.D. Bol, M.R. Clover, R.M. DeVries, N.J. DiGiacomo, J.S. Kapustinsky, P.L. McGaughey, G.R. Smith, W.E. Sondheim, M. Buenerd, J. Chauvin, D. Lebrun, P. Martin, and J.C. Dousse, PS 187-a good statistics study of antiproton interactions with nuclei: preliminary results, Inst. Phys. Conf. Ser. No. 73: Section 3, Paper presented at VII Eur. Symp. Antiproton Interactions, Durham 1984.
14. B.G. Harvey, Introduction to Nuclear Physics and Chemistry, second edition, Prentice-Hall, Inc., Englewood Cliffs, New Jersey, 1969, p.310.
15. M.R. Clover, private communication.
16. K. Nakamura, J. Chiba, T. Fujii, H. Iwasaki, T. Kageyama, S. Kuribayashi, T. Sumiyoshi, T. Takeda, H. Ikeda, and Y. Takada, Absorption and Forward Scattering of Antiprotons by C, Al, and Cu Nuclei in the Region 470-880 MeV/c, Phys. Rev. Lett. 52, 731 (1984).
17. J.H. Kim and H. Toki, Theoretical Study of Pion Spectra after Antiproton Annihilation at Rest in Nuclei, Prog. Theor., Phys. 72 (Progress Letters), 1050 (1984).
18. H. Mes and J. Hébert, Monte-Carlo Calculation of Antiprotons Annihilating (at rest) in Nuclei, Using the Isobar Model, Canadian Journal of Physics 49, 441 (1971).

# BIBLIOGRAPHY

- L.E. Agnew, T. Elioff, W.B. Fowler, R.L. Lander, W.M. Powell, E. Segré, H.M. Steiner, H.S. White, C. Wiegand, and T. Ypsilantis, Antiproton Interactions in Hydrogen and Carbon below 200 MeV, Physical Review 118, 1371 (1960).
- R. Alrichs, O. Dumbrajs, H. Pilkuhn, and H.G. Schlaile, Antiprotonic Helium and Lithium with One or Two Electrons, Z. Phys. A. - Atoms and Nuclei 306, 297 (1982).
- E. Amaldi, G. Baroni, G. Belletini, C. Castagnoli, M. Ferro-Luzzi, and A. Manfredini, Study of Antiprotons with Emulsion Technique; Nuovo Cimento 14, 977 (1959).
- J. Amoros, J.M. Bolta, and F. Senent, Aniquilacion de Antiprotones en Reposo Sobre Nucleos de Emulsion Ionografica, Anales de Fisica 71, 103 (1975).
- "Antinucleon-Nucleon Interactions", Proceedings of the Third European Symposium, Stockholm, July 9-13, 1976, Ed. G. Ekspong and S. Nilsson, Pergamon Press, Oxford, 1977.
- A. Apostolakis, G. Briggs, N.A. Khan, and J.V. Major, The Annihilation of Antiprotons on Protons and Neutrons, Il Nuovo Cimento 37, 4080 (1965).
- G. Backenstoss, M. Hasinoff, P. Pavlopoulos, J. Repond, L. Tauscher, D. Tröster, P. Blum, R. Guigas, H. Koch, M. Meyer, H. Poth, U. Raichz, B. Richter, L. Adiels, I. Bergström, K. Fransson, A. Kerek, M. Suffert, and K. Zioutas, Proton - Antiproton Annihilation at rest into  $\pi^0\omega$ ,  $\pi^0\eta$ ,  $\pi^0\gamma$ ,  $\pi^0\pi^0$  and  $\pi^0\eta'$ , Annual Report on Nuclear Physics Activities, July 1, 1982 - June 30, 1983, Ed. H. Gemmeke, F. Käppeler, Ch. Weddigen, Kernforschungszentrum, Karlsruhe, p. 117 (1983).
- F. Balestra, M.P. Busa, L. Busso, L. Ferrero, D. Pamzneri, G. Piragino, F. Tosello, C. Guaraldo, A. Maggiora, Yu. A. Batusov, I.V. Falomkin, G.B. Pontecorvo, M.G. Sapozhnikov, G. Bendiscioli, V. Filippini, G. Fumagalli, C. Marciano, A. Rotondi, A. Zenoni, E. Lodi Rizzini, M. Vascon, and G. Zanella, Inelastic Interaction of Antiprotons with  $^4\text{He}$  Nuclei between 200 and 600 MeV/c, CERN-EP/85-92, June 12, 1985, European Organization for Nuclear Research, Geneva (preprint, submitted to Physics Letters).
- C. Baltay, N. Barash, P. Franzini, N. Gelfand, L. Kirsch, G. Lütjens, J.C. Severiens, J. Steinberger, D. Tycko, and D. Zanello, Test of Charge-Conjugation Invariance in  $\bar{p}$ -p Annihilations at Rest, Physical Review Letters 15, 591 (1965).
- C. Baltay, N. Barash, P. Franzini, N. Gelfand, L. Kirsch, G. Lütjens, D. Miller, J.C. Severiens, J. Steinberger, T.H. Tan, D. Tycko, D. Zanello, R. Goldberg and R.J. Plano, Annihilations of Antiprotons in Hydrogen at Rest into Two Mesons, Phys. Rev. Lett. 15, 532 (1965); Eratum: Phys. Rev. Lett. 15, 597 (1965).



C. Baltay, D. Franzini, N. Gelfand, G. Lütjens, J.C. Severiens, J. Steinberger, D. Tycko and D. Zanello, Annihilation of Antiprotons in Hydrogen at Rest. II. Analysis of the Annihilation into Three Pions, Physical Review 140B, 1039 (1965).

C. Baltay, P. Franzini, G. Lütjens, J.C. Severiens, J. Steinberger, D. Tycko, and D. Zanello, Annihilation of Antiprotons in Hydrogen at Rest. III. The Reactions  $\bar{p} + p \rightarrow \omega^0 + \pi^+ + \pi^-$  and  $\bar{p} + p \rightarrow \omega^0 + \rho^0$ , Phys. Rev. 140B, 1042 (1965).

C. Baltay, P. Franzini, G. Lütjens, J.C. Severiens, D. Tycko, and D. Zanello, Annihilations of Antiprotons at Rest in Hydrogen. V. Multipion Annihilations, Physical Review 145, 1103 (1966).

N. Barash, P. Franzini, L. Kirsch, D. Miller, J. Steinberger, T.H. Tan, R. Plano, and P. Yeager, Antiproton Annihilation in Hydrogen at Rest. I. Reaction  $\bar{p} + p \rightarrow K + \bar{K} + \pi$ , Physical Review 139, 1659 (1965).

N. Barash, L. Kirsch, D. Miller and T.H. Tan, Annihilations of Antiprotons at Rest in hydrogen. IV.  $\bar{p} p \rightarrow \bar{K} K \pi \pi$ , Phys. Rev. 145, 1095 (1966).

W.H. Barkas, R.W. Birge, W.W. Chupp, A.G. Ekspong, G. Goldhaber, S. Goldhaber, H.H. Heckman, D.H. Perkins, J. Sandweiss, E. Segre, F.M. Smith, D.H. Stark, L. Van Rossum, E. Amaldi, G. Baroni, C. Castagnoli, C. Franzinetti, and A. Manfredini, Antiproton-Nucleon Annihilation Process (Antiproton Collaboration Experiment), Physical Review 105, 1037 (1957).

C.J. Batty, E. Friedman, J. Lichtenstadt, Optical Potentials for Low Energy Antiproton-Nucleus Interactions, Physics Letters 142B, 241 (1984).

Yu. A. Batusov, I.V. Falomkin, G.B. Pontecorvo, M.G. Sapozhnikov, C. Guaraldo, A. Maggiora, M. Vascom, G. Zanella, G. Bendiscioli, V. Filippini, E. Lodi Rizzini, A. Rotondi, A. Zenoni, F. Balestra, M.P. Bussa, L. Busso, L. Ferrero, G. Gervino, D. Panzieri, G. Piragino, and F. Tosello (Dubna-Frascati-Padora-Pavia-Torino Collaboration), Restriction on Amount of Antimatter in the Early Universe from  $\bar{p} - {}^4\text{He}$  Reaction Data, Lettere al Nuova Cimento 41, 223 (1984).

A. Bettini, M. Cresti, A. Grigoletto, S. Limentani, A. Loria, L. Peruzzo and R. Santangelo, Spin and Parity of the 1320 MeV pn Resonance, Il Nuovo Cimento 38 1495 (1965).

A. Bettini, M. Cresti, S. Limentani, L. Peruzzo, R. Santangelo, S. Sartori, L. Bertanza, A. Bigi, R. Carrara, R. Casali, and P. Lariccia, Annihilation Into Pions of the pn System from Antiprotons at Rest in Deuterium, Il Nuovo Cimento 47A, 770 (1967).

L. Bracci, G. Fiorentini, and O. Pitzura, Protonium Formation in Flight, Physics Letters 85B, 280 (1979).

R. Brou, Sections Efficaces d'annihilations Antiprotons-Noyaux, Il Nuovo Cimento 13A, N.1, 257 (1972).

M. Cahay, J. Cugnon, P. Jasselette and J. Vandermeulen, Antiproton Annihilation Inside Nuclei, Physics Letters 115B, 7 (1982).

M. Cahay, J. Cugnon and J. Vandermeulen, Low-Energy Antiproton Annihilation in Nuclei, Nuclear Physics A393, 237 (1983).

O. Chamberlain, G. Goldhaber, L. Jauneau, T. Kalogeropoulos, E. Segré, and R. Silberberg, Antiproton-Nucleon Annihilation Process. II, Physical Review 113, 1615 (1959).

W. Chinowsky and G. Kojoian, Antiproton Annihilation at Rest in Deuterium, Il Nuovo Cimento 43, 3275 (1966).

M.R. Clover, R.M. DeVries, N.J. DiGiacomo, and Y. Yariv, Low Energy Antiproton-Nucleus Interactions, Phys. Rev. C 26, 2138 (1982).

G.T. Condo, T. Handler, H.O. Cohn,  $\Lambda^0$  Production from Low Energy Antiproton Annihilation in Complex Nuclei, Physical Review 29, 1531 (1984).

M. Cresti, A. Grigoletto, S. Limentani, A. Loria, L. Peruzzo, R. Santangelo, G.B. Chadwick, W.T. Davies, M. Derrick, C.J.B. Hawkins, P.M.D. Gray, J.H. Mulvey, P.B. Jones, D. Radojicic, and C.A. Wilkinson, Protonium Annihilation into Pions, Sienna International Conference on Elementary Particles 1, 263 (1963).

J. Cugnon and J. Vandermeulen, A Model of Antiproton Annihilation in Nuclei Implied by Two Nucleons, Inst. Phys. Conf. Ser. No. 73: Section 8, Paper presented at VII Eur. Symp. Antiproton Interactions, Durham 1984, p. 499.

J. Cugnon and J. Vandermeulen, Possibility of Unusual Antiproton Annihilation on Nuclei, Physics Letters 146B, 16 (1984).

N.J. DiGiacomo and M.R. Clover, The Significance of Positive Kaon Production in Low-Energy  $\bar{p}$ -Nucleus Annihilations, J. Phys. G: Nucl. Phys. 10, L119 (1984).

A.G. Ekspong and B.E. Ronne, Antiproton Interaction Cross Sections, Il Nuovo Cimento 13, 1599 (1959).

A.G. Ekspong, A. Frisk, S. Nilsson and B.E. Ronne, Antiproton Annihilations in Complex Nuclei, Nuclear Physics 22, 353 (1961).

J.E. Enstrom, T. Ferbel, P.F. Slattery, B.L. Werner, Z.G.T. Guiragossian, Y. Surni, and T. Yoshida (Particle Data Group),  $\bar{N}N$  and  $\bar{N}D$  Interactions - a Compilation, LBL-58, May 1972, Lawrence Berkeley National Laboratory, Berkeley, California.

I.V. Falomkin, F. Nichitia, and G. Piragino, One  $\pi^-$  and two  $\pi^-$  Emission in Antiproton Annihilation on d and  $^4\text{He}$ , Lettere Al Nuovo Cimento 38, 211 (1983).

R.L. Forward, Antimatter Propulsion, Journal of the British Interplanetary Society 35, 391 (1982).

R.L. Forward, Alternate Propulsion Energy Sources, AFRPL TR-83-067, December 1983, Air Force Rocket Propulsion Laboratory, Edwards Air Force Base, California 93523, pp. 1-1 to 1-24.

R.L. Forward, Antiproton Annihilation Propulsion, AIAA-84-1482, AIAA/SAE/ASME 20th Joint Propulsion Conference, Cincinnati, 11-13 June 1984, AIAA, 1633 Broadway, New York 10019.

R.L. Forward, Recommendation for an Air Force Rocket Propulsion Laboratory Near-Term In-House Research Project on Antiproton Propulsion, AFRPL TR-84-079, February 1985, Air Force Rocket Propulsion Laboratory, Edwards Air Force Base, California 93523-5000.

R.L. Forward, Antiproton Annihilation Propulsion, AFRPL-TR-85-034, September 1985, Air Force Rocket Propulsion Laboratory, Edwards Air Force Base, California 93523-5000.

D. Garreta, P. Birien, G. Bruge, A. Chaumeaux, D.M. Drake, S. Janouin, D. LeGrand, M.C. Mallet-LeMaire, B. Mayer, J. Pain, J.C. Peng, M. Berrada, J.P. Bocquet, E. Monnard, J. Mougey, P. Perrin, E. Aslanides, O. Bing, A. Erell, J. Lichtenstadt, A.I. Yavin, Scattering of Antiprotons from Carbon at 46.8 MeV, Physics Letters 135B, 266 (1984).

U. Gastaldi,  $\bar{p}p$  Experiments at Very Low Energy Using Cooled Antiprotons, Proceedings of the Fourth European Antiproton Conference, Barr, 25-30 June, 1978, Ed. A. Fridman, Centre National de la Recherche Scientifique, Paris, France, Vol. 2, p. 607 (1979)

A. Gsponer and J. Hurni, Antimatter Induced Fusion and Thermonuclear Explosions, Independent Scientific Research Institute, 15, Charles-Galland, CH-1206, Geneva (ISRI 85-05, 27 August 1985, preprint).

N. Horwitz, D. Miller, J. Murray, and R. Tripp, Low-Energy Antiproton Interactions in Hydrogen and Deuterium, Physical Review 115, 472 (1959).

A.S. Iljinov, V.I. Nazarov and S.E. Chigrinov, Nuclear Absorption of Stopped Antiprotons: Multipion-nucleus Interactions, Nuclear Physics A382, 378 (1982).

C. Ishii, Leading Particle Effect on Antiproton Annihilation in Neon, Physical Review C25, 1652 (1982).

M. Jezabek and K. Zalewski, High Energy Antiproton Annihilation on Nuclei and Multiple Scattering Models, Report No. 1061/PH, Kraków, Czerwiec 1979, Institute of Nuclear Physics, Cracow, Poland.

M. Jezabek and K. Zalewski, High Energy Antiproton Annihilation on Nuclei and Multiple Scattering Models, Acta Physica Polonica 1311, 425 (1980).

T.E. Kalogeropoulos, L. Gray, A. Nandy, and J. Roy, Antiproton-nucleon Annihilation into Three Pions, Physical Review D 24, 1959 (1981).

J.H. Kim and H. Toki, Theoretical Study of Pion Spectra after Antiproton Annihilation at Rest in Nuclei, Prog. Theor. Phys. 72, 1050 (1984).

M. Kohno and W. Weise, Quark Model Description of Low Energy Proton-Antiproton Annihilation into Two Mesons, Physics Letters 152B, 303 (1985).

W. Kolos, D.L. Morgan, Jr., D.M. Schrader, and L. Wolniewicz, Hydrogen-Antihydrogen Interactions, Phys. Rev. A 11, 1792 (1975).

L.A. Kondratyuk and M.Zh. Schmatikov, Interaction of Low and Intermediate-Energy Antiprotons with  $^3\text{He}$  and  $^3\text{H}$  Nuclei, Sov. J. Nucl. Phys. 38(2), 216 (1983).

A.D. Mackellar, G.R. Satchler, and C.Y. Wong, An Exploratory Study of Antiproton-Nucleus Scattering, Z. Phys. A-Atoms and Nuclei 316, 35 (1984).

P.F. Massier, The Need for Expanded Exploration of Matter-Antimatter Annihilation for Propulsion Application, Journal of the British Interplanetary Society 35, 387 (1982).

P.L. McGaughey, K.D. Bol, M.R. Clover, R.M. DeVries, N.J. DiGiacomo, J.S. Kapustinsky, G.R. Smith, J.W. Sunier, W.E. Sondheim, Y.Yariv, M. Buenerd, J. Chauvin, D. Lebrun, P. Martin, and J.C. Dousse, Antiproton Annihilations in Nuclei, in "Hadronic Probes and Nuclear Interactions," AIP Conference Proceedings, No. 133, Ed. J.R. Comfort, W.R. Gibbs, B.G. Ritchie, HIPANI Conference, Tempe, Arizona, March 10-14, 1985, American Institute of Physics, New York, 1985, p.41.

P.L. McGaughey, M.R. Clover, and N.J. DiGiacoma,, Antiproton Annihilations in Nuclei May Not Produce High Nuclear Temperatures, LA-UR-85-1266 (Revised), Los Alamos National Laboratory, Los Alamos, New Mexico 87545 (1985) (preprint, submitted to Physics Letters B).

H. Mes, I. Ahmad, J. Hébert, Resonance in the Nucleus, Prog. Theor. Phys. 35, 565 (1966).

H. Mes and J. Hébert, Monte Carlo Calculation of Antiprotons Annihilating (at rest) in Nuclei, Using the Isobar Model, Canadian Journal of Physics 49, 441 (1971).

D.H. Miller, The Elementary Particles with Strong Interactions, Part 2: The Meson Systems, High Energy Physics 4, 111 (1969).

D.L. Morgan, Jr., Concepts for the Design of Antimatter Annihilation Rocket, Journal of the British Interplanetary Society 35, 405 (1982).

D.L. Morgan, Jr., and V.W. Hughes, Atomic Processes Involved in Matter-Antimatter Annihilation, Phys. Rev. D 2, 1389 (1970).

D.L. Morgan, Jr., and V.W. Hughes, Atom-Antiatom Interactions, Phys. Rev. A 7, 1811 (1973).

K. Nakamura, J. Chiba, T. Fujii, H. Iwasaki, T. Kageyama, S. Kuribayashi, T. Sumiyoshi, T. Takeda, H. Ikeda, and Y. Takada, Absorption and Forward Scattering of Antiprotons by C, Al, and Cu Nuclei in the Region 470-880 MeV/c, Physical Review Letters 52, 731 (1984).

V.I. Nazaruk, Antiproton Annihilation on  $^4\text{He}$ , Physics Letters 155B, 323 (1985).

S.F. Orfanidis and V. Rittenberg, Nucleon-Antinucleon Annihilation into Pions, Nuclear Physics B59, 570 (1973).

C. Pajares and A.V. Ramallo, Intermediate Annihilation States in Antiproton-Nucleus Scattering, Physics Letters 107B, 238 (1981).

PANIC Particles and Nuclei, Tenth International Conference, Book of Abstracts, Vol. I and II, Heidelberg, July 30 - August 3, 1984, Ed. F. Güttner, B. Povh, G. zu Putlitz.

P. Pavlopoulos, G. Backenstoss, P. Blüm, K. Fransson, R. Guigas, N. Hassler, M. Izycki, H. Koch, A. Nilsson, H. Poth, M. Suffert, L. Tauscher, and K. Zioutas, Possible Evidence for Narrow Bound States Related to the  $p\bar{p}$  System, Nucleon-Nucleon Interactions - 1977 (Vancouver), AIP Conference Proceedings, No. 41, Ed. D.F. Measday, H.W. Fearing, and A. Strathdee, (Series Ed. H.C. Wolfe), American Institute of Physics, New York, 1978, p. 341 (see also following article by F. Myhrer, Low-Energy Nucleon-Antinucleon Interaction, op. cit., p.357).

E. Pelaquier, Proton-Antiproton Annihilation at Rest and High Energy Reactions, Proceedings of the Fourth European Antiproton Conference, Barr, 25-30 June 1978, Centre National de la Recherche Scientifique, 1979, Vol. 1, p. 51.

Physics at LEAR with Low-Energy Cooled Antiprotons, Ed. U. Gastaldi and R. Klapisch, Plenum Press, New York, 1984.

Physics with Antiprotons at LEAR in the ACOL Era, Proceedings of the Third LEAR Workshop, Tignes, Savoie, France, January 19-26, 1985, Ed. U. Gastaldi, R. Klapisch, J.M. Richard, and J. Tran Thanh Van, editions Frontiers, Yvette, France (1985).

Physikertagung Gemeinsam mit der Fruehjahrstagung der Fachgremien Kern-und Teilchenphysik, Didaktik der Physik, Vakuumtechnik, Relativitaetstheorie, Muenchen (Germany, F.R.), 11-15 March 1985, Verh. Dtsch. Phys. Ges. 20(3), 439 (1985), (Abstracts H3.6 and H3.7).

R. Rüeckl and Ć. Zupančić, Inner Bremsstrahlung in Low Energy Antiproton-Proton Annihilation, Physics Letters 150B, 225 (1985).

J.W. Sumier, K.D. Bol, M.R. Clover, R.M. DeVries, N.J. DiGiacomo, J.S. Kapustinsky, P.L. McGaughey, G.R. Smith, W.E. Sondheim, M. Buenerd, J. Chauvin, D. Lebrun, P. Martin, and J.C. Dousse, PS 187 - a Good Statistics Study of Antiproton Interactions with Nuclei: Preliminary Results, Inst. Phys. Conf. Ser. No. 73: Section 3, p. 195, paper presented at VII Eur. Symp. Antiproton Interactions, Durham 1984.

G. Vulpetti, Antimatter Propulsion for Space Exploration, Preprint IAA-85-491, 36th Congress of the International Astronautical Federation, Stockholm, Sweden, October 7-12, 1985 (Peaceful Space and Global Problems of Mankind), Pergamon Press, Oxford (1985).

M. Wade and V.G. Lind, Ratio of Antiproton Annihilations on Neutrons and Protons in Carbon for Low-Energy and Stopped Antiprotons, Physical Review D 14, 1182 (1976).

END

DTIC

6-86